



CITY UNIVERSITY  
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# Entanglement Measures from Quantum Field Theory Methods in 1+1 Dimensions

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- This talk is mainly based on the papers:

Olivier Blondeau-Fournier, O. C.-A. and Benjamin Doyon,  
*Universal scaling of the logarithmic negativity in massive  
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[arXiv:1508.04026](https://arxiv.org/abs/1508.04026)

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A good review of our work can be found at:

O. C-A and Benjamin Doyon, *Bi-partite entanglement entropy in massive (1+1)-dimensional quantum field theories*, J. Phys. A42 (2009) 504006; [arXiv:0906.2946](#)

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- What provides a good measure of entanglement? [Plenio & Virmani'05]
- The bi-partite entanglement entropy [Bennett et al.'96] and the logarithmic negativity [Peres'96; Eisert'00; Vidal & Werner'01; Plenio'05] are good measures of entanglement according to these properties

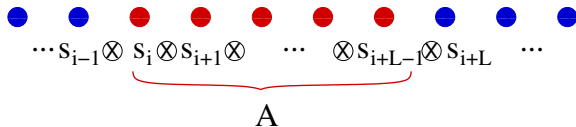
# Bi-partite Entanglement Entropy (EE)

- Let us consider a spin chain of length  $N$ , subdivided into regions  $A$  and  $\bar{A}$  of lengths  $L$  and  $N - L$



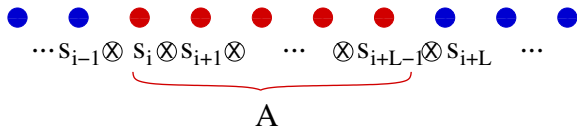
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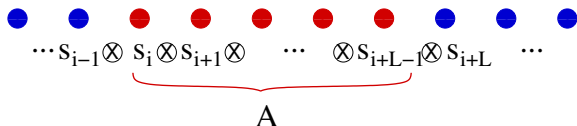
## Von Neumann Entanglement Entropy

$$S_A = -\text{Tr}_A(\rho_A \log(\rho_A)) \quad \text{with} \quad \rho_A = \text{Tr}_{\bar{A}}(|\Psi\rangle\langle\Psi|)$$

$|\Psi\rangle$  ground state and  $\rho_A$  the reduced density matrix.

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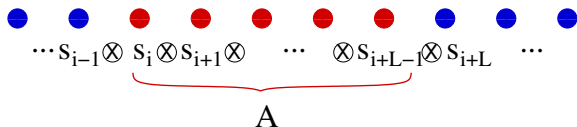
- Other entropies may also be defined such as

## Other Entropies

$$S_A^{\text{Rényi}} = \frac{\log(\text{Tr}_A(\rho_A^n))}{1-n}, \quad S_A^{\text{Tsallis}} = \frac{1 - \text{Tr}_A(\rho_A^n)}{n-1}$$

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## Replica Trick

$$S_A = -\text{Tr}_A(\rho_A \log(\rho_A)) = -\lim_{n \rightarrow 1} \frac{d}{dn} \text{Tr}_A(\rho_A^n)$$

- For general QFTs the “replica trick” naturally leads to the notion of replica theories on multi-sheeted Riemann surfaces  $\Rightarrow$  interpretation of  $\text{Tr}_A(\rho_A^n)$

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## Logarithmic Negativity

$$\mathcal{E} = \log \text{Tr}_{AUB} |\rho_{AUB}^{T_B}| \quad \text{with} \quad \rho_{AUB} = \text{Tr}_C (|\Psi\rangle\langle\Psi|)$$

- Where  $\text{Tr}|\rho|$  represents the sum of the absolute values of the eigenvalues of  $\rho$  and  $T_B$  represents “partial transposition” in sub-system  $B$

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- $|\Psi\rangle$  is the state of the whole system (for pure states)

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- There is also a “replica” approach to the computation of the negativity [Calabrese, Cardy & Tonni'12]:

## Logarithmic Negativity from the Replica Trick

$$\mathcal{E}[n] = \log \text{Tr}_{A \cup B} (\rho_{A \cup B}^{T_B})^n \quad \text{then} \quad \mathcal{E} = \lim_{n \rightarrow 1} \mathcal{E}_e[n]$$

where  $\mathcal{E}_e[n]$  means the function  $\mathcal{E}[n]$  for  $n$  even. This limit requires analytic continuation from  $n$  even to  $n = 1$

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Universal scaling: [Holzhey, Larsen & Wilczek'94; Vidal, Latorre, Rico & Kitaev'03; Calabrese & Cardy'04]:

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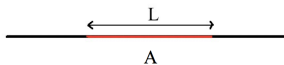
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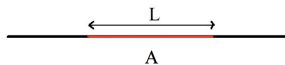
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$$S(L) \sim -\frac{c}{3} \log m\epsilon + U - \frac{a}{8} K_0(2mL)$$

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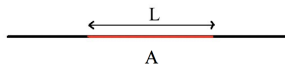
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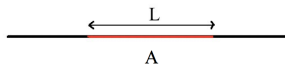
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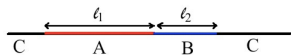
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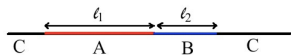
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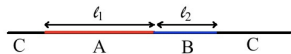
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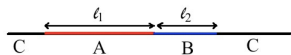
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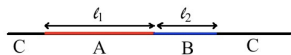
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Near criticality:

Universal saturation and decay: For adjacent regions ( $\ell_1 := \ell, \ell_2 \rightarrow \infty$ ) [Blondeau-Fournier, OC-A & Doyon’16]

$$\mathcal{E}^\perp(\ell) \sim -\frac{c}{4} \log m\epsilon + \mathcal{E}_{\text{sat}} - \frac{2a}{3\sqrt{3}\pi} K_0(\sqrt{3}m\ell)$$

where  $m \propto \xi^{-1}$  is the smallest mass scale in the theory

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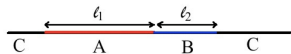
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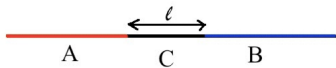
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where  $m \propto \xi^{-1}$  is the smallest mass scale in the theory “ $a$ ” is the number of lightest particles, and  $\mathcal{E}_{\text{sat}}$  is a universal constant.

For semi-infinite non-adjacent regions:



$$\mathcal{E}^{\perp+}(\ell) \sim \frac{a(m\ell)^2}{2\pi^2} \left[ K_0(m\ell)^2 + \frac{K_0(m\ell)K_1(m\ell)}{m\ell} - K_1(m\ell)^2 \right]$$

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 $\Rightarrow$  **Bessel function decay**

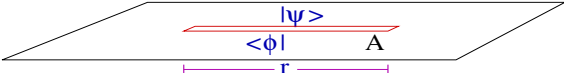
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  - $\Rightarrow$  **Bessel function decay**
- QFT even allows us to access the universal part of the saturation constants (both for the EE and the LN) in the gapped regime
  - $\Rightarrow$  The saturation constants  $U$  and  $\mathcal{E}_{\text{sat}}$  are related to **expectation values** of QFT fields and **structure constants** in CFT

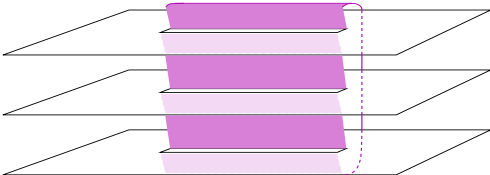
So, how is this done?

# Partition functions on multi-sheeted Riemann surfaces

For integer number  $n$  of replicas, in the scaling limit,  $\text{Tr}_A(\rho_A^n)$  and  $\text{Tr}_{A \cup B}(\rho_{A \cup B}^{T_B})^n$  are partition functions on Riemann surfaces [Callan & Wilczek '94; Holzhey, Larsen & Wilczek '94; Calabrese & Cardy '04; Calabrese, Cardy & Tonni'12]:

$${}_A\langle\phi|\rho_A|\psi\rangle_A \sim \text{Diagram}$$


$$\text{Tr}_A(\rho_A^n) \sim Z_n = \int [d\varphi]_{\mathcal{M}_n} \exp \left[ - \int_{\mathcal{M}_n} d^2x \mathcal{L}[\varphi](x) \right]$$

$$\mathcal{M}_n = \text{Diagram}$$


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- For more complicated measures (e.g. LN or EE of many intervals), even in CFT, it becomes increasingly difficult to obtain analytical expressions
- This is because, in essence, measures of entanglement (in the replica approach) are related to partition functions in Riemann manifolds whose complexity increases with the number of intervals involved

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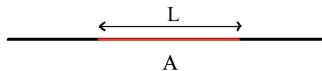
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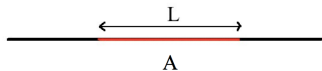


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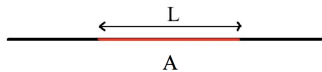
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# EE & LN from Twist Fields

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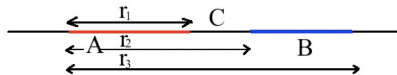
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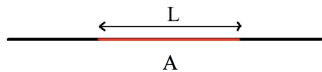


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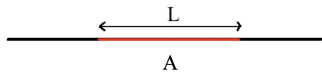
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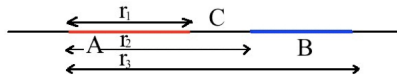
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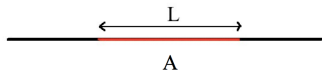
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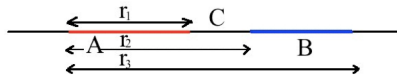
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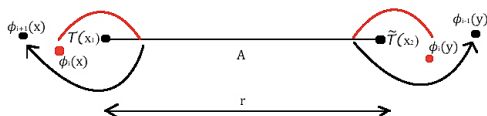
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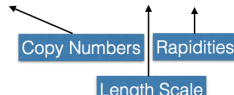
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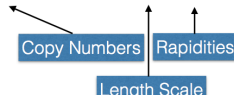
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